## CH<sub>3</sub>O<sup>+</sup> and C<sub>2</sub>H<sub>5</sub>O<sup>+</sup>: High Barriers to Isomerisation and Low Barriers to Symmetry-allowed 1,1-Elimination

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Summary Experimental and theoretical evidence is presented to show that the alkoxy cations  $CH_3O^+$  and  $C_2H_5O^+$  have surprisingly high barriers to 1,2-hydride shifts to the oxygen atom, but have very low barriers to 1,1-elimination of  $H_2$  from the  $\alpha$ -carbon atom.

In earlier work,<sup>1</sup> we have shown that  $CH_2 = OH$  undergoes symmetry-forbidden 1,2-elimination of  $H_2$  in a reaction requiring *ca*. 330 kJ mol<sup>-1</sup>. The high internal energy required to dissociate  $CH_2 = OH$  was shown by deuterium labelling<sup>1,2</sup> to be insufficient to cause any prior isomerisation of the ion. Yet the reported heat of formation of the methoxy cation (CH<sub>3</sub>O<sup>+</sup>, 850 kJ mol<sup>-1</sup>),<sup>3</sup> is only 140 kJ mol<sup>-1</sup> higher than that of CH<sub>2</sub>=OH (710 kJ mol<sup>-1</sup>).<sup>4</sup> These data imply that the barrier to the CH<sub>2</sub>=OH  $\rightarrow$  CH<sub>3</sub>O<sup>+</sup> isomerisation is >330 kJ mol<sup>-1</sup>, and that the barrier for the reverse reaction CH<sub>3</sub>O<sup>+</sup>  $\rightarrow$  CH<sub>2</sub>=OH (occurring through the same channel) is >190 kJ mol<sup>-1</sup>. The latter barrier appears, at first sight, to be a remarkably high one to a symmetry-allowed 1,2-hydride shift which can lead to a

thermodynamically more stable product. In contrast, in the gas phase a primary carbonium ion appears to rearrange to a secondary carbonium ion via a 1,2-hydride shift essentially without activation energy.<sup>5,6</sup>

However, unambiguous experimental evidence for a barrier to the reaction  $CH_3O^+ \rightarrow CH_2=OH$  is available from the work of Hiraoka and Kebarle.<sup>7</sup> Using a pulsed, high-pressure mass spectrometer source, these workers were able to show that in a thermally equilibrated system,  $H_2$  will add to the formyl cation in a reaction which is exothermic by 16 kJ mol<sup>-1</sup>. This exothermicity is calculated from the temperature variation of the equilibrium constant for the reaction, *i.e.* the reaction is readily reversible. The adduct is formulated as  $CH_3O^+$ ,<sup>7</sup> as it clearly must be since 1,2-addition of  $H_2$  to  $HC^+=O$  would require an activation energy of *ca.* 220 kJ mol<sup>-1</sup>, if occurring by microscopic

reversal of 1,2-elimination of  $H_2$  from  $CH_2=OH$ . Moreover, the reaction  $CH_2=OH \rightarrow HC=O + H_2$  could not occur in a system in thermal equilibrium at the temperatures employed (-100 to -165 °C).<sup>7</sup> Thus,  $H_2$  adds 1,1 to the formyl ion, and dissociation of the adduct requires less energy than

isomerisation to CH2=OH.

The product  $CH_3O^+$  from the above  $H_2 + HC^+=O$  reaction has  $\Delta H_f \ 810 \text{ kJ mol}^{-1}$ , and this should be an accurate value.<sup>7</sup> If it is indeed a more accurate value than the appearance potential value (850 kJ mol}^-1) quoted earlier, then the barrier ( $\geq 230 \text{ kJ mol}^{-1}$ ) derived for the  $CH_3O^+ \rightarrow CH_2=OH$  reaction is even higher than the one quoted at the outset of this paper.

We have sought confirmation of the above deductions by

generating  $CH_3O^+$  and  $CH_2=OH$  in the mass spectrometer by ionisation and fragmentation of selected compounds. Since it appears that  $RCH_2O^+$  ions lose  $H_2$  through 1,1elimination with a very small activation energy, precursors of these ions should give very low abundance  $RCH_2O^+$  ions and high abundance RC=O ions. In contrast,  $CH_2=OH$  ions generated from  $RCH_2OH$  compounds should give large m/e31:29 ratios since the m/e 31  $\rightarrow$  29 reaction has a very large activation energy. The experimental results are given in the Table.

TABLE TABLE M/e 31:29 ratios in the mass spectra (70 eV) of selected compounds

| Compound   | Anticipated ion<br>at threshold | m/e 31:29      |
|--|---------------------------------|----------------|
| CH₃CH₂OH   | $CH_2 = OH$                     | 11:1           |
| CF <sub>3</sub> CH <sub>2</sub> OH                 | CH₂=ÓH                          | 11:1           |
| CH <sub>3</sub> CH <sub>2</sub> CH <sub>2</sub> OH | CH <sub>2</sub> =OH             | 15:1           |
| CH <sub>3</sub> ONO <sub>2</sub>                   | $CH_{3}O^{+}$                   | 0.08:1         |
| CH.OCH.  | $CH_{O}^{+}$                    | $0.08 \cdot 1$ |

The results strikingly confirm expectations, both  $CH_3$ - +ONO<sub>2</sub> and  $CH_3OCH_3$  producing extremely abundant HC=O ions. Furthermore, while  $CH_2$ =OH loses  $H_2$  in slow reactions (metastable transitions) with a large kinetic energy release<sup>2</sup> (140 kJ mol<sup>-1</sup>) since the reaction is symmetryforbidden,<sup>1</sup> the  $CH_3O^+ \rightarrow HC=O$  reaction does not give a discernible metastable peak. Presumably this is because the symmetry-allowed reaction (see below) has a small activation energy and a relatively fast rise of the rate constant with energy, such that almost all the dissociations are relatively fast.

Finally, we wished to establish that an acyl ion RC=0, when formed from a compound  $RCH_2OX$ , does arise specifically by 1,1-elimination of  $H_2$  from the oxygenbound carbon even where other processes are possible. Ethyl nitrate produces in its 70 eV mass spectrum peaks at m/e 43 (B) and 45 (A) in the ratio 20:1 (Scheme 1).  $[1,1-^{2}H_{2}]$ -Ethyl nitrate produces under similar conditions peaks at m/e 43 (B) and 47 (C) in the ratio 12:1, with peaks at m/e 44 and 45 being of negligible abundance relative to that at m/e 43. The  $CH_{3}C=0$  ion is therefore produced by

that at m/e 43. The CH<sub>3</sub>C=O ion is therefore produced by a specific 1,1-elimination.

$$\begin{array}{c} -\mathbf{e} & -\mathbf{H}_{2} & +\\ \mathrm{CH}_{3}\mathrm{CH}_{2}\mathrm{ONO}_{2} \xrightarrow{-\mathbf{e}} \mathrm{CH}_{3}\mathrm{CH}_{2}\mathrm{O}^{+} \xrightarrow{-\mathbf{H}_{2}} & +\\ -\mathrm{NO}_{2} & \\ & (\mathbf{A}), \ m/e \ 45 & (\mathbf{B}), \ m/e \ 43 \\ \mathrm{CH}_{3}\mathrm{CD}_{2}\mathrm{ONO}_{2} \xrightarrow{-\mathbf{e}} \mathrm{CH}_{3}\mathrm{CD}_{2}\mathrm{O}^{+} \xrightarrow{-\mathrm{D}_{2}} & +\\ -\mathrm{NO}_{2} & \\ & -\mathrm{NO}_{2} & \\ & (\mathbf{C}), \ m/e \ 47 & (\mathbf{B}), \ m/e \ 43 \\ \mathrm{SCHEME} \ 1 \end{array}$$

The extremely small barrier to the reaction  $CH_3O^+ \rightarrow HC=O + H_2$  is supported by earlier orbital symmetry arguments.<sup>8</sup> 1,1-Elimination from an  $sp^3$  carbon adjacent to a vacant *p*-orbital is a symmetry-allowed process.



It remains to explain the high barrier ( $\ge 230 \text{ kJ mol}^{-1}$ ) to the reaction  $\text{CH}_3\text{O}^+ \rightarrow \text{CH}_2=\text{OH}$ . The hydride shift should occur so that bonding on to the vacant *p*-orbital on oxygen is possible in the transition state (1)  $\rightarrow$  (2) (Scheme 2). However, this results in the formation of  $\text{CH}_2\text{OH}$  in which the lone-pairs of electrons on the oxygen atom are orthogonal to the vacant  $\pi$ -orbital on the carbon atom (3). Thus (3), formed in this manner, lacks all the  $\pi$ -stabilisa-

tion of  $CH_2$ =OH, but suffers all the destabilisation of the +  $CH_2OH$  cation which exists through  $\sigma$ -electron withdrawal by the electronegative oxygen atom; it may plausibly

represent the highest point on the potential surface of the  $CH_3O^+ \rightarrow CH_2=OH$  isomerisation. The heat of formation of (3) may be estimated in the following manner. The difference in heats of formation of  $CH_4$  and  $CH_3^+$  is 1155 kJ mol<sup>-1</sup>;<sup>9</sup> the analogous change performed on methanol  $[CH_3OH \rightarrow (3)]$  should require a similar energy difference if  $\pi$  back-donation to the cationic centre did not occur, nor did

inductive destabilisation by oxygen. Lack of  $\pi$ -donation is precisely the situation which we wish to examine in (3). Inductive destabilisation will however be present in (3); calculations by Pople and his co-workers<sup>10</sup> give an inductive destabilisation of a carbocation of  $42 \text{ kJ} \text{ mol}^{-1}$  when the oxygen atom is separated by one carbon atom from the cationic centre, and 13 kJ mol<sup>-1</sup> when it is separated by two carbon atoms. In view of the rapid decrease of inductive effects due to an increase in the number of intervening  $\sigma$ -bonds, it seems probable that the inductive destabilisation due to a directly bonded oxygen will not be

less than 100 kJ mol<sup>-1</sup>. We therefore conclude that  $\Delta H_{f}$ (3) is  $\ge 1055 \text{ kJ mol}^{-1}$ . Thus, a plausible transition state for the reversible reaction  $CH_2 = OH \rightleftharpoons CH_3 - O$  appears to be energetically inaccessible at the high internal energies necessary to induce the reaction  $CH_2=OH \rightarrow HC=O + H_2$ ; this is in accord with the experimental facts.

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